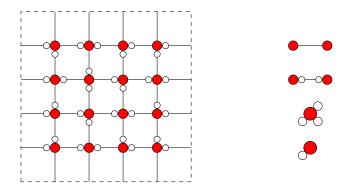
# Computing the entropy of multidimensional subshifts of finite type

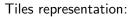
Silvère Gangloff

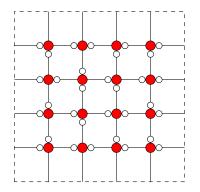
LIP, ENS Lyon

Septembre 11, 2018

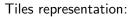
#### Bidimensional ice stable states [Pauling-Lieb]: Forbidden:

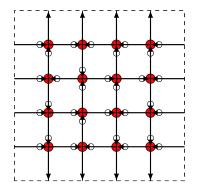




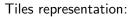


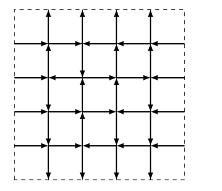






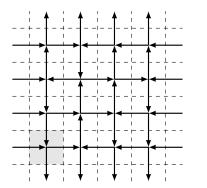






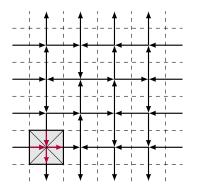


## **Tiles representation:**



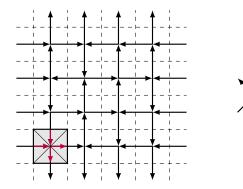


## **Tiles representation:**





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H. Wang Tiles (1960'):



possible:



not possible:



Ex: Hard square shift, ou hard core model.

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Alphabet  $\mathcal{A} = \{0, 1\}$ , forbidden patterns  $\begin{bmatrix} 1\\ 1 \end{bmatrix}$  et  $\boxed{1 \pm 1}$ .

Ex: Hard square shift, ou hard core model.

Entropy: "quantity of possible states".

 $N_n(X)$ : number of observable *n* size squares.



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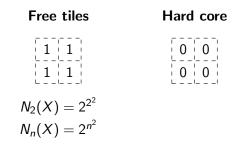


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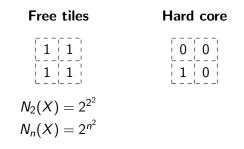
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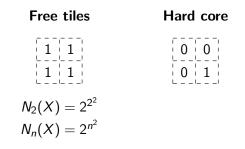


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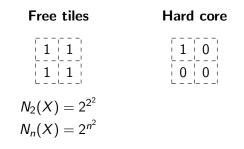
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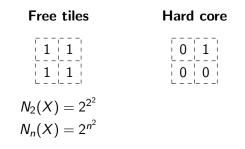
Entropy: "quantity of possible states".



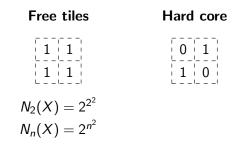
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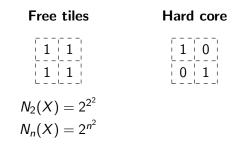
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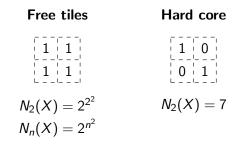
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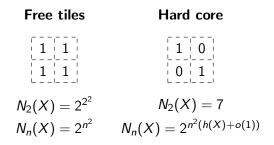
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$$h(X) = \inf_{n \ge 1} \frac{\log_2(N_n(X))}{n^2}$$

Free tiles	Hard core	Ice [Lieb 67]
h=1	$h\geq 1/2$	$h = (4/3)^{3/2}$ ?

$$h(X) = \inf_{n \ge 1} \frac{\log_2(N_n(X))}{n^2} = \inf_n \frac{\log_2(N_n^{loc}(X))}{n^2}$$

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Free tiles	Hard core	Ice [Lieb 67]
h=1	$h\geq 1/2$	$h = (4/3)^{3/2}$ ?

$$n = 1$$
  $\longrightarrow$  Algorithme  $\xrightarrow{h(X)}$ 

$$h(X) = \inf_{n \ge 1} \frac{\log_2(N_n(X))}{n^2} = \inf_n \frac{\log_2(N_n^{loc}(X))}{n^2}$$

Free tiles	Hard core	Ice [Lieb 67]
h=1	$h\geq 1/2$	$h = (4/3)^{3/2}$ ?

$$n = 2$$
  $\longrightarrow$  Algorithme  $\xrightarrow{h(X)}$ 

$$h(X) = \inf_{n \ge 1} \frac{\log_2(N_n(X))}{n^2} = \inf_n \frac{\log_2(N_n^{loc}(X))}{n^2}$$

Free tiles	Hard core	Ice [Lieb 67]
h = 1	$h\geq 1/2$	$h = (4/3)^{3/2}$ ?

$$n = 3$$
  $\longrightarrow$  Algorithme  $\xrightarrow{h(X)}$ 

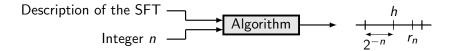
$$h(X) = \inf_{n \ge 1} \frac{\log_2(N_n(X))}{n^2} = \inf_n \frac{\log_2(N_n^{loc}(X))}{n^2}$$

Free tiles	Hard core	Ice [Lieb 67]
h=1	$h\geq 1/2$	$h = (4/3)^{3/2}$ ?

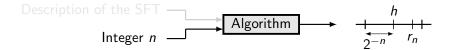
$$n = 4$$
  $\longrightarrow$  Algorithme  $\xrightarrow{h(X)}$ 

## Does there exist a "universal method" to compute the entropy of SFT?

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Implies computability:



The possible values of the bidimensional SFT are the **semi-computable** from above numbers.

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Exemple:

$$\sum_{k\in\mathcal{H}}2^{-k},$$

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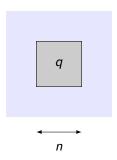
These numbers include non-computable numbers.

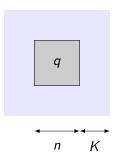
Exemple:

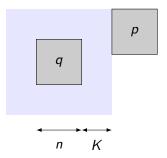
$$\sum_{k\in\mathcal{H}} 2^{-k},$$

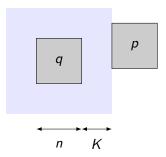
 $\mathcal{H}$ : programs which never stop on n = 0.

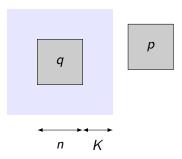






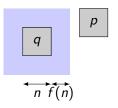




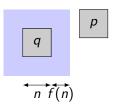


f(n)-block gluing:

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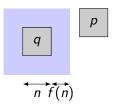


### Théorème (G., Sablik)

**1** The possible values of entropy for linear block gluing (f(n) = O(n), ex: ice) are the non-negative numbers semi-computable from above.

2 If  $f(n) = o(\log(n))$ , entropy is computable.

f(n)-block gluing:



### Théorème (G., Sablik)

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2 If  $f(n) = o(\log(n))$ , entropy is computable.

Answer to problem 9.1 of Hochman and Meyerovitch.

#### Characterisation of a threshold for f?

For **decidable subshifts** (not including all SFT), meaning that  $n \mapsto N_n(X)$  is computable:

## Théorème (G., Hellouin)

Let  $f : \mathbb{N} \to \mathbb{N}$  be non-decreasing.

- **1**  $\sum_{n} \frac{f(n)}{n^2} < +\infty$  with computable speed: the entropy is computable.
- 2  $\sum_{n} \frac{f(n)}{n^2} = +\infty$ : possible values: semi-computable numbers from above.

Existence of a threshold for multidimensional SFT?

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Characterisation of the possible values under the threshold ?

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Characterisation of the possible values under the threshold ?

**General aim:** understand the conditions under which we can compute the entropy of physical pertinent models.